

## Analytical expansions for Fermi–Dirac functions

Alicia Bañuelos, Ricardo Angel Depine, and Roberto Claudio Mancini

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# Analytical expansions for Fermi-Dirac functions

Alicia Bañuelos <sup>a)</sup>

Laboratorio de Física del Plasma, Dirección General de Investigación y Desarrollo, Ministerio de Defensa  
Departamento de Física-Facultad de Ciencias Exactas y Naturales, Universidad de Buenos Aires, 1428 Buenos Aires, Argentina

Ricardo Angel Depine <sup>a)</sup>

Laboratorio de Optica, Departamento de Física-Facultad de Ciencias Exactas y Naturales, Universidad de Buenos Aires, 1428 Buenos Aires, Argentina

Roberto Claudio Mancini <sup>a)</sup>

Laboratorio de Física del Plasma, Dirección General de Investigación y Desarrollo, Ministerio de Defensa  
Departamento de Física-Facultad de Ciencias Exactas y Naturales, Universidad de Buenos Aires, 1428 Buenos Aires, Argentina

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We obtain a fast convergent series expansion for the Fermi-Dirac function  $F_\sigma(\alpha)$  for  $-10 \leq \alpha \leq -1$ . We give values of  $F_\sigma(\alpha)$  for  $\sigma = n + \frac{1}{2}$  ( $n = 0, 1, \dots, 6$ ) with  $\alpha$  in the same range.

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## I. INTRODUCTION

The Fermi-Dirac functions  $F_\sigma(\alpha)$ , where  $\sigma$  is a positive real parameter, is defined for all real numbers  $\alpha$  by

$$F_\sigma(\alpha) = \frac{1}{\Gamma(\sigma)} \int_0^\infty \frac{x^{\sigma-1}}{e^{x+\alpha} + 1} dx.$$

When  $\sigma$  is an integer, this integral may be easily evaluated by a power series; a complete discussion of this case is due to Rhodes.<sup>1</sup> For arbitrary  $\sigma$ , there are several expansions depending on the range of values of  $\alpha$ .<sup>2-4</sup> The calculation of  $F_\sigma(\alpha)$  for  $\alpha < 0$  is needed in many questions of quantum statistical mechanics; for example, to solve the equations of state corresponding to extreme conditions (high pressure and nonzero temperature). Analytical expansions are available in all ranges, except when  $-10 \leq \alpha \leq -1$ . Previous evaluations of  $F_\sigma(\alpha)$  for this range were made by numerical integration<sup>4,5</sup> or by polynomial approximation.<sup>6</sup> In this paper we obtain a fast convergent series expansion of  $F_\sigma(\alpha)$  for  $-10 \leq \alpha \leq -1$ .

## II. SERIES EXPANSION FOR $-10 \leq \alpha \leq -1$

For simplicity, let us define

$$I = I_\sigma(\alpha) = \Gamma(\sigma) F_\sigma(\alpha) = \int_0^\infty \frac{x^{\sigma-1}}{e^{x+\alpha} + 1} dx.$$

Substituting  $y = x + \alpha$  in this integral gives

$$I = \int_{-\alpha}^\infty \frac{(y + |\alpha|)^{\sigma-1}}{e^y + 1} dy.$$

Now  $I$  will be calculated as

$$I = I_1 + I_2 + I_3$$

by dividing the integration interval by the points  $-p$  and  $p$ , where  $0 < p < |\alpha|$ . Another restriction on the values of  $p$  and convenient numerical suggestions will appear later.

## A. Evaluation of $I_1$

First we expand the integrand denominator of

$$I_1 = \int_{-\alpha}^{-p} \frac{(y + |\alpha|)^{\sigma-1}}{e^y + 1} dy \quad (1)$$

in a series of powers of  $e^y$

$$\frac{1}{e^y + 1} = \sum_{n \geq 0} (-1)^n e^{ny}.$$

This series converges uniformly for  $|e^y| < 1$ , that is, for  $y < 0$ . Now, by expanding  $e^{ny}$  at a convenient point  $y_0$ , we obtain

$$e^{ny} = e^{ny_0} \sum_{k \geq 0} \frac{n^k (y - y_0)^k}{k!}.$$

By replacing successively in (1), taking into account the uniform convergence of the series to exchange the order of integrals and summations, it follows that

$$I_1 = \sum_{n \geq 0} (-1)^n e^{ny_0} \sum_{k \geq 0} \frac{n^k}{k!} \times \int_{-\alpha}^{-p} (y + |\alpha|)^{\sigma-1} (y - y_0)^k dy.$$

The integrals involved in this expression may be evaluated using the formula

$$\int (a + bx)^{\sigma-1} x^k dx = \frac{(a + bx)^\sigma}{b^{\sigma+1}} \times \sum_{0 \leq j \leq k} \frac{(-1)^j \binom{k}{j} (a + bx)^{k-j} a^j}{k - j + \sigma}.$$

Thus,

$$I_1 = (|\alpha| - p)^\sigma \sum_{n \geq 0} (-1)^n e^{ny_0} \sum_{k \geq 0} \frac{n^k}{k!} A_k,$$

where

$$A_k = \sum_{0 \leq j \leq k} \frac{(-1)^j \binom{k}{j} (|\alpha| - p)^{k-j} (y_0 + |\alpha|)^j}{k - j + \sigma}.$$

<sup>a)</sup>Fellow of CONICET (Consejo Nacional de Investigaciones Científicas y Técnicas, Argentina).

**B. Evaluation of  $I_2$**

We notice that

$$I_2 = |\alpha|^{\sigma-1} \int_{-p}^p \frac{(y/|\alpha| + 1)^{\sigma-1}}{e^y + 1} dy.$$

Since  $0 < p < |\alpha|$ , the series

$$\left(\frac{y}{|\alpha|} + 1\right)^{\sigma-1} = \sum_{n \geq 0} \binom{\sigma-1}{n} \left(\frac{y}{|\alpha|}\right)^n$$

converges uniformly. The same statement holds<sup>7</sup> for the series

$$\frac{1}{e^y + 1} = \frac{1}{2} + \frac{1}{2} \sum_{k \geq 1} \frac{(1-4^k)B_k y^{2k-1}}{k(2k-1)!}$$

when  $p < \pi$ ,  $B_k$  being the nonzero Bernoulli numbers. Therefore, arguments used in Sec. (A) apply here, yielding

$$I_2 = |\alpha|^{\sigma-1} \sum_{n \geq 1} C_n \left( \frac{1}{\sigma - 2n + 1} + \frac{1}{|\alpha|} \sum_{k \geq 1} \frac{D_k}{2k + 2n - 1} \right),$$

where

$$C_1 = (\sigma - 1)p,$$

$$C_{n+1} = C_n \frac{(\sigma - 2n)(\sigma - 2n - 1)}{(2n + 1)2n} \left(\frac{p}{|\alpha|}\right)^2,$$

$$D_1 = -3B_1 p^2,$$

and

$$D_{k+1} = D_k \frac{(1 - 4^{k+1})B_{k+1} p^2}{(1 - 4^k)B_k 2(k+1)(2k+1)}.$$

**C. Evaluation of  $I_3$**

Recall that

$$I_3 = \int_p^\infty \frac{(y + |\alpha|)^{\sigma-1}}{e^y + 1} dy.$$

Since the expansion of the integrand denominator in a series of powers of  $e^{-y}$

$$\frac{1}{e^y + 1} = - \sum_{n \geq 1} (-1)^n e^{-ny},$$

is uniformly convergent for  $y > 0$ , by exchanging the integral and the summation, with the substitution  $z = n(y + |\alpha|)$  we obtain

$$I_3 = \sum_{n \geq 1} \frac{(-1)^{n+1} e^{n|\alpha|}}{n^\sigma} \int_{n(|\alpha|+p)}^\infty z^{\sigma-1} e^{-z} dz.$$

Thus,  $I_3$  can be expressed in terms of incomplete Gamma functions as

$$I_3 = \sum_{n \geq 1} \frac{(-1)^{n+1} e^{n|\alpha|}}{n^\sigma} \Gamma(\sigma, n(p + |\alpha|)).$$

From our numerical investigations we conclude that, in order to achieve a fast convergence, the values of  $p$  and  $y_0$  may be chosen as follows:

$$p = |\alpha|/2, \text{ if } 1 \leq |\alpha| < 5;$$

$$p = 2.5, \text{ if } 5 \leq |\alpha| < 10$$

$$y_0 = -(|\alpha| + p)/2.$$

As an application, values of  $F_\sigma(\alpha)$  for  $-10 \leq \alpha \leq -1$  and  $\sigma = n + \frac{1}{2}$  ( $n = 0, 1, \dots, 6$ ) were computed with a maximum relative error of  $10^{-5}$ . In particular, we have checked the accuracy of all previously tabulated values.<sup>4,6</sup> In the course of the computation we have made use of Abramowitz's tables<sup>7</sup> for Bernoulli numbers. The corresponding computing program is to be published elsewhere.<sup>8</sup>

TABLE I. Values of  $F_\sigma(\alpha)$  for  $-10 \leq \alpha \leq -1$  and  $\sigma = n + \frac{1}{2}$  with  $n = 0, 1, \dots, 6$ .

| ALPHA | F1/2        | F3/2        | F5/2        | F7/2        | F9/2        | F11/2       | F13/2       |
|-------|-------------|-------------|-------------|-------------|-------------|-------------|-------------|
| -1.0  | 1.027050 00 | 1.515650 00 | 2.002250 00 | 2.294840 00 | 2.478670 00 | 2.587180 00 | 2.648300 00 |
| -1.1  | 1.071660 00 | 1.600580 00 | 2.165020 00 | 2.503120 00 | 2.718420 00 | 2.846840 00 | 2.919800 00 |
| -1.2  | 1.116280 00 | 1.709970 00 | 2.338520 00 | 2.728200 00 | 2.979860 00 | 3.131580 00 | 3.218490 00 |
| -1.3  | 1.160820 00 | 1.803840 00 | 2.523170 00 | 2.971190 00 | 3.264670 00 | 3.443600 00 | 3.547010 00 |
| -1.4  | 1.205190 00 | 2.002130 00 | 2.719430 00 | 3.233230 00 | 3.574720 00 | 3.785360 00 | 3.908200 00 |
| -1.5  | 1.249320 00 | 2.144870 00 | 2.927750 00 | 3.515470 00 | 3.911980 00 | 4.159460 00 | 4.309160 00 |
| -1.6  | 1.293150 00 | 2.271990 00 | 3.148550 00 | 3.819180 00 | 4.278550 00 | 4.568720 00 | 4.741270 00 |
| -1.7  | 1.336600 00 | 2.403480 00 | 3.382290 00 | 4.145630 00 | 4.676590 00 | 5.016200 00 | 5.220190 00 |
| -1.8  | 1.379640 00 | 2.539250 00 | 3.629400 00 | 4.496100 00 | 5.108450 00 | 5.505160 00 | 5.745900 00 |
| -1.9  | 1.422220 00 | 2.679390 00 | 3.890300 00 | 4.871950 00 | 5.576640 00 | 6.039130 00 | 6.322730 00 |
| -2.0  | 1.464290 00 | 2.823730 00 | 4.165410 00 | 5.274620 00 | 6.083740 00 | 6.621810 00 | 6.955340 00 |
| -2.1  | 1.505840 00 | 2.972230 00 | 4.455170 00 | 5.705520 00 | 6.632530 00 | 7.257260 00 | 7.648820 00 |
| -2.2  | 1.546830 00 | 3.124870 00 | 4.760000 00 | 6.166170 00 | 7.225860 00 | 7.949790 00 | 8.408660 00 |
| -2.3  | 1.587250 00 | 3.281580 00 | 5.080290 00 | 6.658050 00 | 7.856830 00 | 8.704000 00 | 9.240840 00 |
| -2.4  | 1.627080 00 | 3.442300 00 | 5.416450 00 | 7.182750 00 | 8.558550 00 | 9.524830 00 | 1.015170 01 |
| -2.5  | 1.666310 00 | 3.606980 00 | 5.768880 00 | 7.741880 00 | 9.304480 00 | 1.041750 01 | 1.114820 01 |
| -2.6  | 1.704940 00 | 3.775550 00 | 6.137970 00 | 8.337070 00 | 1.010810 01 | 1.138770 01 | 1.223780 01 |
| -2.7  | 1.742980 00 | 3.947950 00 | 6.524110 00 | 8.970030 00 | 1.097320 01 | 1.244120 01 | 1.342850 01 |

(continued)

TABLE I. (Continued).

|      |             |             |             |             |             |             |             |
|------|-------------|-------------|-------------|-------------|-------------|-------------|-------------|
| -2.8 | 1.780410 00 | 4.124120 00 | 6.927690 00 | 9.642480 00 | 1.190340 01 | 1.358450 01 | 1.472910 01 |
| -2.9 | 1.817240 00 | 4.304010 00 | 7.349060 00 | 1.035620 01 | 1.290300 01 | 1.482420 01 | 1.614800 01 |
| -3.0 | 1.853490 00 | 4.487550 00 | 7.788610 00 | 1.111290 01 | 1.397610 01 | 1.616750 01 | 1.769730 01 |
| -3.1 | 1.889150 00 | 4.674690 00 | 8.246690 00 | 1.191450 01 | 1.512710 01 | 1.762200 01 | 1.938500 01 |
| -3.2 | 1.924240 00 | 4.865360 00 | 8.723670 00 | 1.276290 01 | 1.636060 01 | 1.919570 01 | 2.122570 01 |
| -3.3 | 1.958770 00 | 5.059520 00 | 9.215880 00 | 1.365990 01 | 1.768130 01 | 2.089700 01 | 2.322920 01 |
| -3.4 | 1.992750 00 | 5.257100 00 | 9.735680 00 | 1.460750 01 | 1.909420 01 | 2.273500 01 | 2.540970 01 |
| -3.5 | 2.026190 00 | 5.458050 00 | 1.027140 01 | 1.560770 01 | 2.060460 01 | 2.471910 01 | 2.778110 01 |
| -3.6 | 2.059120 00 | 5.662320 00 | 1.082740 01 | 1.666250 01 | 2.221760 01 | 2.685930 01 | 3.035870 01 |
| -3.7 | 2.091530 00 | 5.869860 00 | 1.140410 01 | 1.777390 01 | 2.393890 01 | 2.916620 01 | 3.315850 01 |
| -3.8 | 2.123450 00 | 6.080610 00 | 1.200150 01 | 1.894390 01 | 2.577430 01 | 3.165090 01 | 3.619790 01 |
| -3.9 | 2.154890 00 | 6.294530 00 | 1.262020 01 | 2.017490 01 | 2.772970 01 | 3.432510 01 | 3.949510 01 |
| -4.0 | 2.185870 00 | 6.511570 00 | 1.326050 01 | 2.146670 01 | 2.991140 01 | 3.720110 01 | 4.306960 01 |
| -4.1 | 2.216390 00 | 6.731690 00 | 1.392260 01 | 2.282770 01 | 3.202570 01 | 4.029180 01 | 4.694240 01 |
| -4.2 | 2.246480 00 | 6.954840 00 | 1.460690 01 | 2.425400 01 | 3.437920 01 | 4.361090 01 | 5.113560 01 |
| -4.3 | 2.276140 00 | 7.180970 00 | 1.531370 01 | 2.574980 01 | 3.687880 01 | 4.717250 01 | 5.567270 01 |
| -4.4 | 2.305390 00 | 7.410050 00 | 1.604320 01 | 2.731750 01 | 3.953150 01 | 5.099170 01 | 6.057870 01 |
| -4.5 | 2.334250 00 | 7.642040 00 | 1.679580 01 | 2.895920 01 | 4.234470 01 | 5.508420 01 | 6.588010 01 |
| -4.6 | 2.362710 00 | 7.876890 00 | 1.757170 01 | 3.067740 01 | 4.532590 01 | 5.946630 01 | 7.160520 01 |
| -4.7 | 2.390810 00 | 8.114570 00 | 1.837130 01 | 3.247440 01 | 4.848280 01 | 6.415520 01 | 7.778370 01 |
| -4.8 | 2.418550 00 | 8.355040 00 | 1.919470 01 | 3.435250 01 | 5.182350 01 | 6.916900 01 | 8.444710 01 |
| -4.9 | 2.445940 00 | 8.598270 00 | 2.004240 01 | 3.631410 01 | 5.535610 01 | 7.452630 01 | 9.162890 01 |
| -5.0 | 2.472990 00 | 8.844220 00 | 2.091450 01 | 3.836180 01 | 5.908920 01 | 8.024690 01 | 9.936440 01 |
| -5.1 | 2.499710 00 | 9.092850 00 | 2.181130 01 | 4.049780 01 | 6.303140 01 | 8.635120 01 | 1.076910 02 |
| -5.2 | 2.526120 00 | 9.344150 00 | 2.273310 01 | 4.272480 01 | 6.719180 01 | 9.286050 01 | 1.166400 02 |
| -5.3 | 2.552220 00 | 9.598070 00 | 2.368020 01 | 4.504530 01 | 7.157950 01 | 9.979710 01 | 1.262770 02 |
| -5.4 | 2.578030 00 | 9.854580 00 | 2.465280 01 | 4.746170 01 | 7.620410 01 | 1.071840 02 | 1.366230 02 |
| -5.5 | 2.603550 00 | 1.011370 01 | 2.565120 01 | 4.997670 01 | 8.107510 01 | 1.150460 02 | 1.477300 02 |
| -5.6 | 2.628790 00 | 1.037530 01 | 2.667560 01 | 5.259290 01 | 8.620280 01 | 1.234080 02 | 1.596480 02 |
| -5.7 | 2.653760 00 | 1.063940 01 | 2.772640 01 | 5.531270 01 | 9.159720 01 | 1.322960 02 | 1.724290 02 |
| -5.8 | 2.678480 00 | 1.090600 01 | 2.880360 01 | 5.813900 01 | 9.726880 01 | 1.417370 02 | 1.861260 02 |
| -5.9 | 2.702940 00 | 1.117510 01 | 2.990760 01 | 6.107440 01 | 1.032290 02 | 1.517590 02 | 2.007960 02 |
| -6.0 | 2.727150 00 | 1.144660 01 | 3.103870 01 | 6.412140 01 | 1.094870 02 | 1.623920 02 | 2.164900 02 |
| -6.1 | 2.751120 00 | 1.172050 01 | 3.219700 01 | 6.728300 01 | 1.150570 02 | 1.736670 02 | 2.332960 02 |
| -6.2 | 2.774870 00 | 1.199680 01 | 3.338290 01 | 7.056180 01 | 1.229480 02 | 1.856140 02 | 2.512540 02 |
| -6.3 | 2.798390 00 | 1.227550 01 | 3.459650 01 | 7.396050 01 | 1.301730 02 | 1.982670 02 | 2.704420 02 |
| -6.4 | 2.821690 00 | 1.255650 01 | 3.583810 01 | 7.748200 01 | 1.377440 02 | 2.116600 02 | 2.909320 02 |
| -6.5 | 2.844770 00 | 1.283980 01 | 3.710790 01 | 8.112910 01 | 1.456740 02 | 2.258280 02 | 3.128000 02 |
| -6.6 | 2.867660 00 | 1.312540 01 | 3.840610 01 | 8.490450 01 | 1.539740 02 | 2.408080 02 | 3.361250 02 |
| -6.7 | 2.890340 00 | 1.341330 01 | 3.973300 01 | 8.881120 01 | 1.626590 02 | 2.566360 02 | 3.609900 02 |
| -6.8 | 2.912820 00 | 1.370350 01 | 4.108880 01 | 9.285210 01 | 1.717410 02 | 2.733530 02 | 3.874820 02 |
| -6.9 | 2.935120 00 | 1.399590 01 | 4.247380 01 | 9.703000 01 | 1.812340 02 | 2.909980 02 | 4.156910 02 |
| -7.0 | 2.957230 00 | 1.429050 01 | 4.388810 01 | 1.013480 02 | 1.911520 02 | 3.096130 02 | 4.457140 02 |
| -7.1 | 2.979160 00 | 1.458730 01 | 4.533200 01 | 1.058090 02 | 2.015080 02 | 3.292430 02 | 4.776480 02 |
| -7.2 | 3.000920 00 | 1.488630 01 | 4.680560 01 | 1.104150 02 | 2.123180 02 | 3.499300 02 | 5.115970 02 |
| -7.3 | 3.022500 00 | 1.518750 01 | 4.830930 01 | 1.151710 02 | 2.235960 02 | 3.717220 02 | 5.476710 02 |
| -7.4 | 3.043920 00 | 1.549080 01 | 4.984320 01 | 1.200780 02 | 2.353570 02 | 3.946660 02 | 5.859800 02 |
| -7.5 | 3.065170 00 | 1.579630 01 | 5.140760 01 | 1.251400 02 | 2.476170 02 | 4.188100 02 | 6.266440 02 |
| -7.6 | 3.086270 00 | 1.610390 01 | 5.300250 01 | 1.303610 02 | 2.603910 02 | 4.442060 02 | 6.697840 02 |
| -7.7 | 3.107200 00 | 1.641350 01 | 5.462840 01 | 1.357420 02 | 2.736950 02 | 4.709060 02 | 7.155280 02 |
| -7.8 | 3.127990 00 | 1.672530 01 | 5.628530 01 | 1.412870 02 | 2.875450 02 | 4.989630 02 | 7.640100 02 |
| -7.9 | 3.148630 00 | 1.703910 01 | 5.797350 01 | 1.470000 02 | 3.019580 02 | 5.284340 02 | 8.153680 02 |
| -8.0 | 3.169120 00 | 1.735500 01 | 5.969320 01 | 1.528830 02 | 3.169500 02 | 5.593740 02 | 8.697460 02 |
| -8.1 | 3.189480 00 | 1.767300 01 | 6.144460 01 | 1.589400 02 | 3.325400 02 | 5.918440 02 | 9.272940 02 |
| -8.2 | 3.209690 00 | 1.799290 01 | 6.322790 01 | 1.651730 02 | 3.487440 02 | 6.259030 02 | 9.881680 02 |
| -8.3 | 3.229770 00 | 1.831490 01 | 6.504320 01 | 1.715860 02 | 3.655810 02 | 6.616130 02 | 1.052530 03 |
| -8.4 | 3.249710 00 | 1.863890 01 | 6.689090 01 | 1.781830 02 | 3.830670 02 | 6.990400 02 | 1.120550 03 |
| -8.5 | 3.269530 00 | 1.896480 01 | 6.877110 01 | 1.849660 02 | 4.012230 02 | 7.382490 02 | 1.192400 03 |
| -8.6 | 3.289220 00 | 1.929280 01 | 7.068390 01 | 1.919380 02 | 4.200670 02 | 7.793080 02 | 1.268260 03 |
| -8.7 | 3.308780 00 | 1.962270 01 | 7.262970 01 | 1.991040 02 | 4.396170 02 | 8.222860 02 | 1.348320 03 |
| -8.8 | 3.328220 00 | 1.995450 01 | 7.460850 01 | 2.064650 02 | 4.598940 02 | 8.672560 02 | 1.432780 03 |
| -8.9 | 3.347540 00 | 2.028830 01 | 7.662070 01 | 2.140260 02 | 4.809170 02 | 9.142900 02 | 1.521840 03 |
| -9.0 | 3.366750 00 | 2.062400 01 | 7.866630 01 | 2.217900 02 | 5.027060 02 | 9.634650 02 | 1.615710 03 |

(continued)

Cont. from Table I.

|       |             |             |             |             |             |             |             |
|-------|-------------|-------------|-------------|-------------|-------------|-------------|-------------|
| -9.1  | 3.38583D 00 | 2.09617D 01 | 8.07455D 01 | 2.29761D 02 | 5.25282D 02 | 1.01486D 03 | 1.71461D 03 |
| -9.2  | 3.40481D 00 | 2.13012D 01 | 8.28587D 01 | 2.37941D 02 | 5.48665D 02 | 1.06855D 03 | 1.81876D 03 |
| -9.3  | 3.42367D 00 | 2.16426D 01 | 8.50058D 01 | 2.46334D 02 | 5.72877D 02 | 1.12462D 03 | 1.92840D 03 |
| -9.4  | 3.44243D 00 | 2.19859D 01 | 8.71873D 01 | 2.54943D 02 | 5.97939D 02 | 1.18315D 03 | 2.04377D 03 |
| -9.5  | 3.46108D 00 | 2.23311D 01 | 8.94031D 01 | 2.63772D 02 | 6.23873D 02 | 1.24423D 03 | 2.16511D 03 |
| -9.6  | 3.47962D 00 | 2.26781D 01 | 9.16535D 01 | 2.72825D 02 | 6.50701D 02 | 1.30796D 03 | 2.29270D 03 |
| -9.7  | 3.49806D 00 | 2.30270D 01 | 9.39388D 01 | 2.82104D 02 | 6.78446D 02 | 1.37441D 03 | 2.42680D 03 |
| -9.8  | 3.51640D 00 | 2.33777D 01 | 9.62590D 01 | 2.91614D 02 | 7.07130D 02 | 1.44368D 03 | 2.56768D 03 |
| -9.9  | 3.53464D 00 | 2.37303D 01 | 9.86144D 01 | 3.01357D 02 | 7.36776D 02 | 1.51586D 03 | 2.71563D 03 |
| -10.0 | 3.55278D 00 | 2.40847D 01 | 1.01005D 02 | 3.11338D 02 | 7.67409D 02 | 1.59106D 03 | 2.87095D 03 |

<sup>1</sup>P. Rhodes, Proc. Roy. Soc. A **204**, 396 (1950).

<sup>2</sup>A. Sommerfeld, Z. Phys. **47**, 1 (1928).

<sup>3</sup>L. Nordheim, "Muller-Pomille's Lehrbuch der Physik" **4** (1934) Braunschweig, as quoted in Ref. 4.

<sup>4</sup>J. McDougall and E. C. Stoner, Philos. Trans. A **237**, 67 (1938).

<sup>5</sup>A. C. Beer, M. N. Chase, and P. F. Choquard, Helv. Phys. Acta **28**, 529

(1955).

<sup>6</sup>R. Latter, Phys. Rev. **99**, 1854 (1955).

<sup>7</sup>M. Abramowitz and I. Stegun, *Handbook of Mathematical Functions, with Formulas, Graphs and Mathematical Tables* (Dover, New York, 1972).

<sup>8</sup>A. Bañuelos, R. A. Depine, and R. C. Mancini (To be published in Comput. Phys. Commun.)